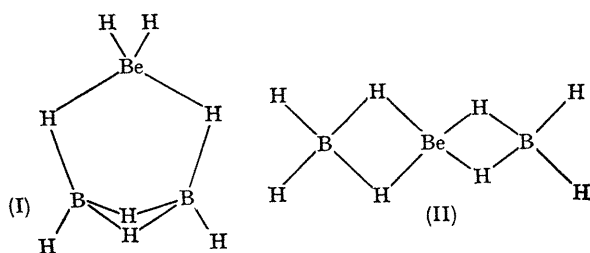


## The Dipole Moment and Infrared Spectrum of Beryllium Borohydride

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PARTICULAR interest in  $\text{BeB}_2\text{H}_8$  stems from a recent electron diffraction study<sup>1</sup> which suggests that the structure has a  $C_{2v}$  configuration (I). Earlier electron diffraction data<sup>2</sup> and i.r. results<sup>3</sup> were interpreted in terms of a  $D_{2d}$  structure (II). Subsequent theoretical calculations on both structures also favour the  $D_{2d}$  configuration.<sup>4,5</sup>



We have re-examined the i.r. spectrum and have measured the dipole moment of  $\text{BeB}_2\text{H}_8$ . We report the symmetry of the molecule.

The dipole moment measurement was frustrated by decomposition of gaseous  $\text{BeB}_2\text{H}_8$  above  $35^\circ$  in the dielectric cell, possibly due to a catalytic effect by the metal electrodes. Although this prohibited a temperature-dependence study of the molar polarization  $P$ , we measured an absolute value of  $P = 105 \pm 40$  c.c./mole at  $25^\circ$ . Coupled with our estimate of the electronic contribution to the molar polarization of  $20 \pm 5$  c.c./mole, we obtain a value of

$2.1 \pm 0.5$  D for the dipole moment. Clearly structure (II) is unacceptable.

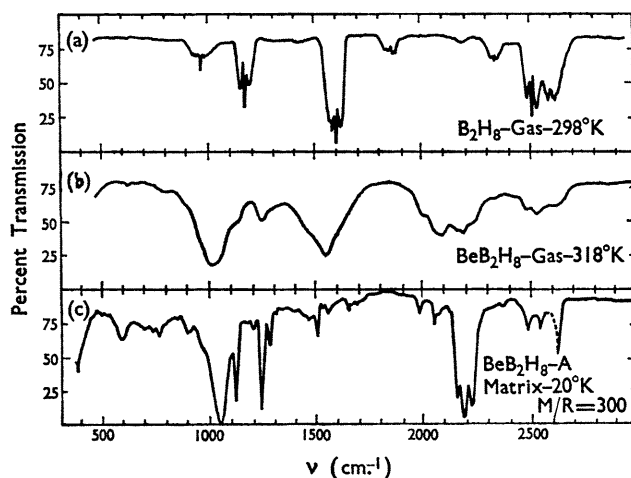


FIGURE. I.r. spectra of  $\text{B}_2\text{H}_6$  and  $\text{BeB}_2\text{H}_8$ .

Further evidence for this is given by the i.r. spectrum (Figure). Identical gas-phase spectra were obtained with inert  $\text{BaF}_2$ ,  $\text{AgCl}_2$ , and quartz windows, thus eliminating the window-reaction problem.<sup>3</sup>

A  $D_{2d}$  configuration of  $\text{BeB}_2\text{H}_8$  would be a symmetric

top molecule, as is  $B_2H_6$  to a good approximation. Hence the parallel vibrational modes should show characteristic  $PQR$  rotational structure. Such structure is apparent in the  $B_2H_6$  spectrum at 2525, 1606, and 1177  $cm^{-1}$ ; however, it is absent in the  $BeB_2H_8$  spectrum, a result best explained if the molecule has  $C_{2v}$  symmetry and hence is an asymmetric top. In addition, calculations of the rotational constants of both structures of  $BeB_2H_8$  lead us to expect resolvable  $Q$ -branches in the perpendicular modes for (II) but not for (I). Such branches are resolved for  $B_2H_6$  (e.g., mode at 2612  $cm^{-1}$ ) but the highest available resolution (0.5  $cm^{-1}$ ) reveals no structure in the  $BeB_2H_8$  absorptions.

The sharp matrix spectra indicate that none of the overlapping gas-phase absorptions are rotational side bands.

Thus, the large number of i.r. absorptions (more than 20 in the matrix) is also an argument against (II) since this structure would have 12 active modes vs. 23 for (I).

The striking drop in intensity of the 1550 and 2075  $cm^{-1}$  gas-phase bands upon matrix isolation is ascribed to an effect such as has been reported for other molecules [ $AlF_3$ ,  $NF_3$ ,  $NO_2$ ,  $CuF_2$ , and  $Cu(NO_3)_2$ ],<sup>6,7</sup> which orient with the molecular plane parallel to the substrate. By symmetry, a similar orientation of the  $BeB_2$  plane in  $BeB_2H_8$  necessitates assignment of one of the 2075 and 1550  $cm^{-1}$  absorptions to the  $BeH$  antisymmetric stretch. The long  $BeH$  bond-length reported for (I) (1.65 Å) suggest assignment of the lower frequency.

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